

$$\begin{aligned}
 u &= u', \quad v = v' \sin \varphi + w \cos \varphi, \quad w = -v \cos \varphi + w \sin \varphi, \\
 \frac{du}{dt} &= \frac{du'}{dt}, \quad \frac{dv}{dt} = \frac{dv'}{dt} \sin \varphi + \frac{dw}{dt} \cos \varphi, \quad \frac{dw}{dt} = -\frac{dv'}{dt} \cos \varphi + \frac{dw'}{dt} \sin \varphi, \\
 u' &= u, \quad v' = v \sin \varphi - w \cos \varphi, \quad w' = v \cos \varphi + w \sin \varphi, \quad (6) \\
 \frac{du'}{dt} &= \frac{du}{dt}, \quad \frac{dv'}{dt} = \frac{dv}{dt} \sin \varphi - \frac{dw}{dt} \cos \varphi, \quad \frac{dw'}{dt} = \frac{dv}{dt} \cos \varphi + \frac{dw}{dt} \sin \varphi.
 \end{aligned}$$

Now at the same time,  $\omega_x = \omega_y = 0, \omega_z = \omega$ , i.e., according to (1),

$$\begin{aligned}
 \frac{du'}{dt} - 2\omega v' + \frac{1}{\rho} \frac{\partial p}{\partial x'} + \frac{\partial \Phi}{\partial x'} &= 0, \\
 \frac{dv'}{dt} + 2\omega u' + \frac{1}{\rho} \frac{\partial p}{\partial y'} + \frac{\partial \Phi}{\partial y'} &= 0, \\
 \frac{dw'}{dt} + \frac{1}{\rho} \frac{\partial p}{\partial z'} + \frac{\partial \Phi}{\partial z'} &= 0.
 \end{aligned} \quad (16)$$

Setting, as in (1)

$$Q = RT_0 \ln \frac{p}{p_0} + \Phi, \quad c^2 = \frac{c_p}{c_v} RT_0, \quad \vartheta = \frac{T - T_0}{T_0},$$

we have

$$\begin{aligned}
 u'_t - 2\omega v' + Q_{x'} &= -[\vartheta(Q_{x'} - \Phi_{x'}) + u'u'_x + v'u'_y + w'u'_z] \equiv F'_1, \\
 v'_t + 2\omega u' + Q_{y'} &= -[\vartheta(Q_{y'} - \Phi_{y'}) + u'v'_x + v'v'_y + w'v'_z] \equiv F'_2, \\
 w'_t + Q_{z'} &= -[\vartheta(Q_{z'} - \Phi_{z'}) + u'w'_x + v'w'_y + w'w'_z] \equiv F'_3, \\
 \frac{1}{c^2} Q_t + u'_x + v'_y + w'_z &= \\
 = -\frac{1}{c^2} [u'(Q_{x'} - \Phi_{x'}) + v'(Q_{y'} - \Phi_{y'}) + w'(Q_{z'} - \Phi_{z'})] &\equiv F'_4
 \end{aligned} \quad (7)$$

Considering here that the magnitude  $1/c^2$  is negligibly slight, and excluding  $u', v',$  and  $w'$  from (7), we get for  $Q$  the equation

$$\begin{vmatrix} \frac{\partial}{\partial t}, & -2\omega, & 0, & \frac{\partial}{\partial x} \\ 2\omega, & \frac{\partial}{\partial t}, & 0, & \frac{\partial}{\partial y} \\ 0 & 0, & \frac{\partial}{\partial t}, & \frac{\partial}{\partial z} \\ \frac{\partial}{\partial x}, & \frac{\partial}{\partial y}, & \frac{\partial}{\partial z}, & 0 \end{vmatrix} Q = \begin{vmatrix} \frac{\partial}{\partial t}, & -2\omega, & 0, & F'_1 \\ 2\omega, & \frac{\partial}{\partial t}, & 0, & F'_2 \\ 0 & 0, & \frac{\partial}{\partial t}, & F'_3 \\ \frac{\partial}{\partial x}, & \frac{\partial}{\partial y}, & \frac{\partial}{\partial z}, & F'_4 \end{vmatrix} \equiv F^{(Q)}. \quad (8)$$

Opening the determinant on the left side we have

$$\left(\frac{\partial^2}{\partial t^2} \Delta_1 + 4\omega^2 \frac{\partial^2}{\partial z^2}\right) Q = -F^{(Q)}, \quad (9)$$

where

$$\Delta_1 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}.$$

The solution of (9) for limitless space has the form [1]

$$\begin{aligned}
 Q &= -\frac{1}{4\pi} \iiint_{-\infty}^{+\infty} \left[ \Delta_1 Q_1 \frac{1}{R_1} J_0 \left( \frac{2\omega r' t}{R_1} \right) + \Delta_1 Q_2 \frac{1}{R_1} \int_0^t J_0 \left( \frac{2\omega r' \tau}{R_1} \right) d\tau \right] d\bar{x}' d\bar{y}' d\bar{z}' + \\
 &+ \iiint_{-\infty}^{+\infty} \left[ -(F'_{1x'} + F'_{2y'} + F'_{3z'}) \frac{1}{R_1} \frac{\partial}{\partial t} J_0 \left( \frac{2\omega r' (t - \tau)}{R_1} \right) + 2\omega (F'_{1y'} - F'_{2x'}) \times \right.
 \end{aligned}$$